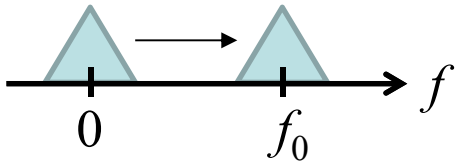
# VIII. Motions on the Time-Frequency Distribution

Fourier spectrum 為 1-D form，只有二種可能的運動或變形：

Modulation $e^{j2\pi f_0 t}x(t) \xrightarrow{FT} X(f-f_0)$

Scaling $x(t/a) \xrightarrow{FT} |a|X(af)$




Time-frequency analysis 為 2-D，在 2-D 平面上有多種可能的運動或變形

(1) Horizontal shifting　　(2) Vertical shifting

(3) Dilation　　(4) Shearing

(5) Generalized Shearing　　(6) Rotation

(7) Twisting

# 8-1 Basic Motions

(1) Horizontal Shifting

$$x(t-t_0) \rightarrow S_x(t-t_0, f)e^{-j2\pi f t_0} \quad \text{,STFT, Gabor}$$
$$\rightarrow W_x(t-t_0, f) \quad \text{,Wigner}$$

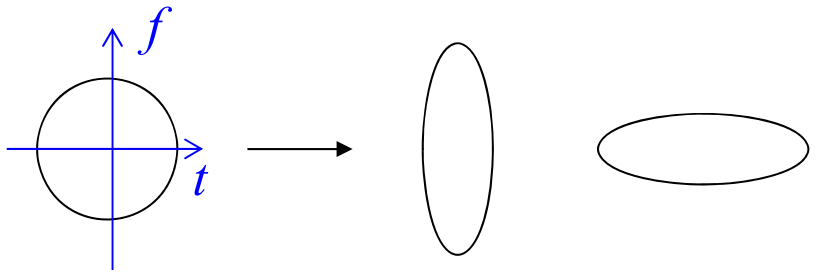
(2) Vertical Shifting

$$e^{j2\pi f_0 t}x(t) \rightarrow S_x(t, f-f_0) \quad \text{,STFT,Gabor}$$
$$\rightarrow W_x(t, f-f_0) \quad \text{,Wigner}$$

## (3) Dilation (scaling)

$$\frac{1}{\sqrt{|a|}} x(\frac{t}{a}) \rightarrow \approx S_x(\frac{t}{a}, af) \text{ ,STFT,Gabor}$$

$$\rightarrow W_x(\frac{t}{a}, af) \text{ ,WDF}$$

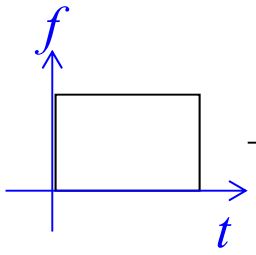


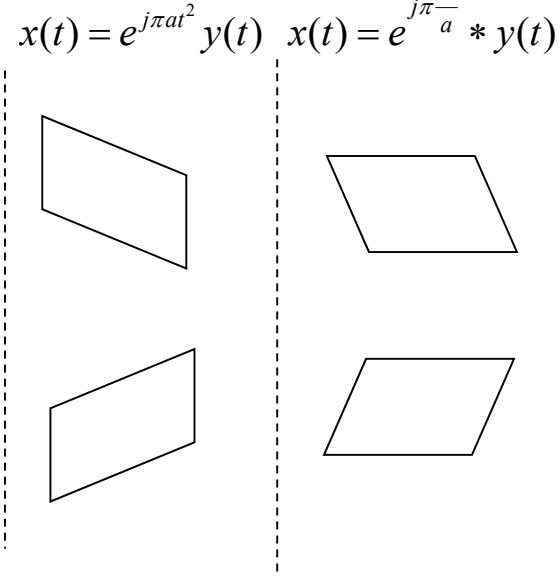
## (4) Shearing

$$x(t) = e^{j\pi a t^2} y(t)$$

$$S_x(t,f) \approx S_y(t, f - at) \text{ ,STFT,Gabor}$$

$$W_x(t,f) = W_y(t, f - at) \text{ ,WDF}$$




$$x(t) = e^{j\pi \frac{t^2}{a}} * y(t) \quad (* \text{ means convolution})$$

$$S_x(t,f) \approx S_y(t - af, f) \text{ , STFT,Gabor}$$

$$W_x(t,f) = W_y(t - af, f) \text{ , WDF}$$

**(Proof):** When $x(t) = e^{j\pi at^2} y(t)$,

$$
\begin{aligned}
W_x(t,f) &= \int_{-\infty}^{\infty} x(t+\tau/2)\, x^*(t-\tau/2)\, e^{-j2\pi\tau f} \cdot d\tau \\
&= \int_{-\infty}^{\infty} e^{j\pi a(t+\tau/2)^2} e^{-j\pi a(t-\tau/2)^2} y(t+\tau/2)\, y^*(t-\tau/2)\, e^{-j2\pi\tau f} d\tau \\
&= \int_{-\infty}^{\infty} e^{j2\pi at\tau} y(t+\tau/2)\, y^*(t-\tau/2)\, e^{-j2\pi\tau f} d\tau \\
&= \int_{-\infty}^{\infty} y(t+\tau/2)\, y^*(t-\tau/2)\, e^{-j2\pi\tau(f-at)} d\tau \\
&= W_y(t, f-at)
\end{aligned}
$$

## (5) Generalized Shearing

$x(t) = e^{j\phi(t)} y(t)$ (What is its effect?)

$$\phi(t) = \sum_{k=0}^{n} a_k t^k$$

$S_x(t, f) \cong S_y(t, f -$ ) ,STFT,Gabor

$W_x(t, f) \cong W_y(t, f -$ ) ,WDF

J. J. Ding, S. C. Pei, and T. Y. Ko, "Higher order modulation and the efficient sampling algorithm for time variant signal," *European Signal Processing Conference*, pp. 2143-2147, Bucharest, Romania, Aug. 2012.

J. J. Ding and C. H. Lee, "Noise removing for time-variant vocal signal by generalized modulation," *APSIPA ASC*, pp. 1-10, Kaohsiung, Taiwan, Oct. 2013

Q:

If $x(t) = h(t) * y(t)$ where $h(t) = IFT\left( \exp\left( j\sum_{k=0}^{n} a_k f^k \right) \right)$

then

$$S_x(t,f) \cong S_y(t + \frac{1}{2\pi}\sum_{k=1}^{n} k a_k f^{k-1}, \quad f) \text{ ,STFT,Gabor}$$

$$W_x(t,f) \cong W_y(t + \frac{1}{2\pi}\sum_{k=1}^{n} k a_k f^{k-1}, \quad f) \text{ ,WDF}$$

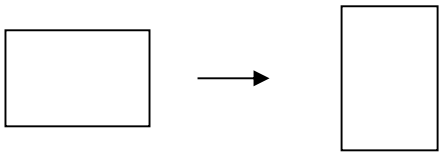
# 8-2 Rotation by $\pi/2$: Fourier Transform

$$X(f) = FT(x(t))$$

$$|S_X(t,f)| \approx |S_x(-f,t)| \quad \text{,STFT}$$

$$G_X(t,f) = G_x(-f,t)e^{-j2\pi ft} \quad \text{,Gabor}$$

$$W_X(t,f) = W_x(-f,t) \quad \text{,WDF}$$



(clockwise rotation by 90°)

Strictly speaking, the rec-STFT have no rotation property.

For Gabor transforms, if

$$G_x(t,f)=\int_{-\infty}^{\infty} e^{-\pi(\tau-t)^2} e^{-j2\pi f\tau} x(\tau)\,d\tau \;,$$

$$G_X(t,f)=\int_{-\infty}^{\infty} e^{-\pi(\tau-t)^2} e^{-j2\pi f\tau} X(\tau)\,d\tau \qquad X(f)=FT\left[x(t)\right]=\int_{-\infty}^{\infty} x(t) e^{-j2\pi f t}\,dt$$

then $G_X(t,f)=G_x(-f,t)e^{-j2\pi t f}$

(clockwise rotation by 90° for amplitude)

(Proof): $G_X(t,f)=\int_{-\infty}^{\infty} e^{-\pi(\tau-t)^2} e^{-j2\pi f\tau}\int_{-\infty}^{\infty} x(u) e^{-j2\pi\tau u}\,du\,d\tau$

$$=\int_{-\infty}^{\infty} x(u) e^{-\pi(\tau-t)^2}\left(\int_{-\infty}^{\infty} e^{-j2\pi\tau(f+u)}\,d\tau\right)du$$

$$=\int_{-\infty}^{\infty} x(u)\left(\int_{-\infty}^{\infty} e^{-\pi(\tau-t)^2} e^{-j2\pi\tau(f+u)}\,d\tau\right)du=\int_{-\infty}^{\infty} x(u)\left(FT\left(e^{-\pi(\tau-t)^2}\right)\Big|_{f\to f+u}\right)du$$

Since $\quad FT\left(e^{-\pi\tau^2}\right)=e^{-\pi f^2}, \quad FT\left(e^{-\pi(\tau-t)^2}\right)=e^{-j2\pi tf}e^{-\pi f^2}$

$$G_X(t,f)=\int_{-\infty}^{\infty} x(u) e^{-j2\pi t(f+u)} e^{-\pi(f+u)^2}\,du$$

$$=e^{-j2\pi tf}\int_{-\infty}^{\infty} x(u) e^{-j2\pi tu} e^{-\pi(u-(-f))^2}\,du = G_x(-f,t)e^{-j2\pi t f}$$

If we define the Gabor transform as

$$G_x(t,f) = e^{j\pi f t}\int_{-\infty}^{\infty} e^{-\pi(\tau-t)^2} e^{-j2\pi f\tau} x(\tau)\,d\tau \quad ,$$

and $G_X(t,f) = e^{j\pi f t}\int_{-\infty}^{\infty} e^{-\pi(\tau-t)^2} e^{-j2\pi f\tau} X(\tau)\,d\tau$

then $G_X(t,f) = G_x(-f,t)$

If $W_x(t,f)=\int_{-\infty}^{\infty} x(t+\tau/2)\cdot x^*(t-\tau/2)e^{-j2\pi\tau f}\cdot d\tau$ is the WDF of $x(t)$,

$W_X(t,f)=\int_{-\infty}^{\infty} X(t+\tau/2)\cdot X^*(t-\tau/2)e^{-j2\pi\tau f}\cdot d\tau$ is the WDF of $X(f)$,

then $W_X(t,f)=W_x(-f,t)$

(clockwise rotation by 90°)

Question: Do other time-frequency distributions also have the similar rotation property?

- If $X(f) = IFT[x(t)] = \int_{-\infty}^{\infty} x(t) e^{j2\pi f t} dt$ , then

$$W_X(t,f) = W_x(f,-t), \qquad G_X(t,f) = G_x(f,-t) e^{j2\pi t f}$$

(counterclockwise rotation by 90°).

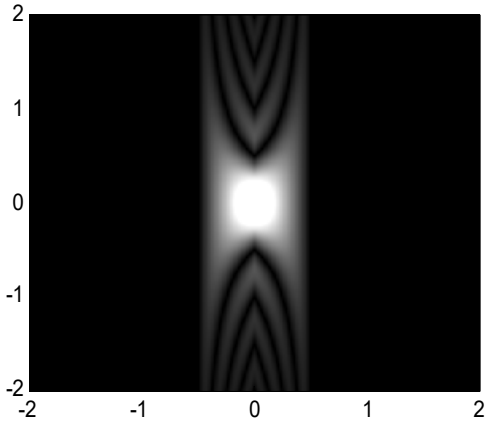
- If $X(f) = x(-t)$ , then

$$W_X(t,f) = W_x(-t,-f) \ , \qquad G_X(t,f) = G_x(-t,-f).$$

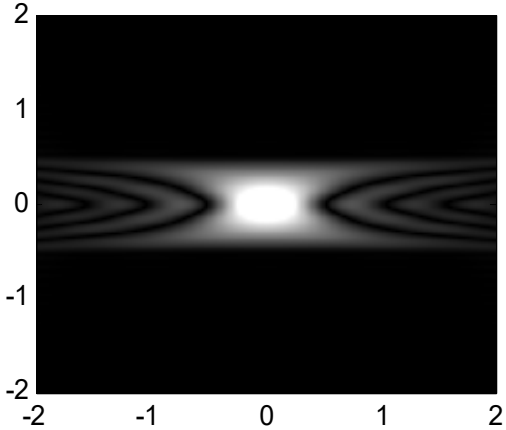
(rotation by 180°).

Examples: $x(t) = \Pi(t)$, $X(f) = FT[x(t)] = \mathrm{sinc}(f)$.

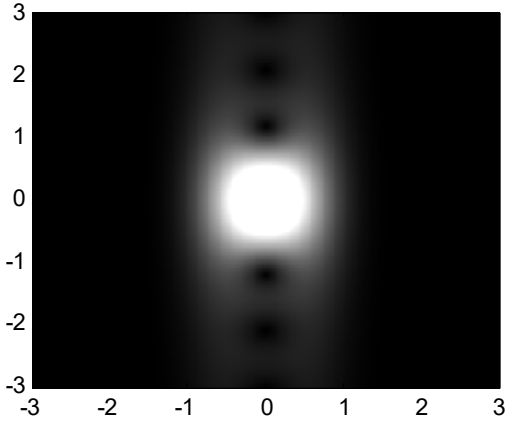
WDF of $\Pi(t)$



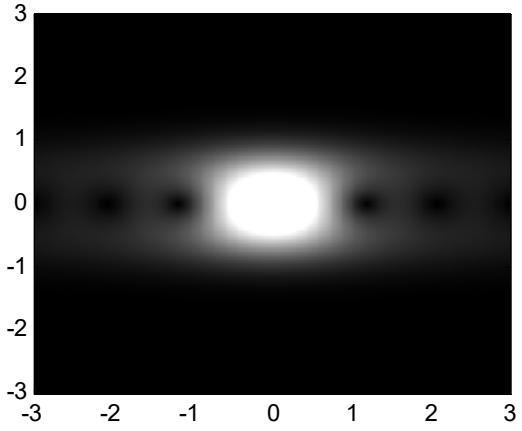

WDF of sinc( $t$ )




Gabor transform of $\Pi(t)$




Gabor transform of sinc( $t$ )



If a function is an eigenfunction of the Fourier transform,

$$\int_{-\infty}^{\infty} e^{-j2\pi f t} x(t)\,dt = \lambda x(f) \qquad \lambda = 1, -j, -1, j$$

then its WDF and Gabor transform have the property of

$$W_x(t,f) = W_x(f,-t) \qquad |G_x(t,f)| = |G_x(f,-t)|$$

(轉了 90º之後，和原來還是一樣)

Example: Gaussian function

$$\exp\left(-\pi t^2\right)$$

## Hermite-Gaussian function

$$\phi_m(t) = \exp\left(-\pi t^2\right) H_m(t)$$

Hermite polynomials: $H_m(t) = C_m e^{2\pi t^2} \frac{d^m}{dt^m} e^{-2\pi t^2}$ , $C_m$ is some constant,

$$H_0(t) = 1 \qquad H_1(t) = t \qquad H_2(t) = 4\pi t^2 - 1$$

$$H_3(t) = 4\pi t^3 - 3t \qquad H_4(t) = 16\pi^2 t^4 - 24\pi t^2 + 3$$

$$\int_{-\infty}^{\infty} e^{-2\pi t^2} H_m(t) H_n(t) = D_m \delta_{m,n} \quad , D_m \text{ is some constant,}$$

$\delta_{m,n} = 1$ when $m = n$, $\delta_{m,n} = 0$ otherwise.

[Ref] M. R. Spiegel, *Mathematical Handbook of Formulas and Tables*, McGraw-Hill, 1990.

Hermite-Gaussian functions are eigenfunctions of the Fourier transform

$$\int_{-\infty}^{\infty} \phi_m(t) e^{-j2\pi f t} dt = (-j)^m \phi_m(f)$$

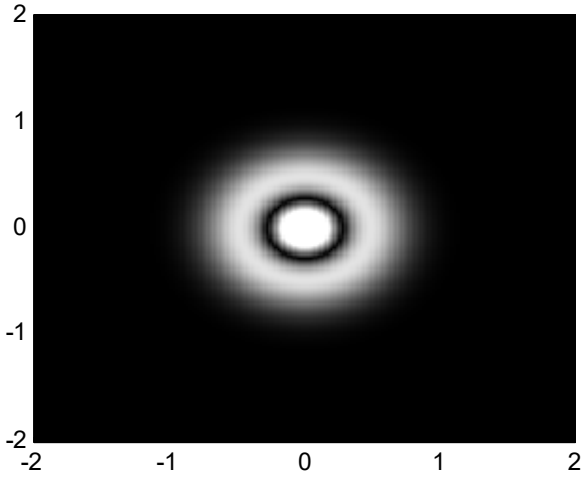
Any eigenfunction of the Fourier transform can be expressed as the form of

$$k(t) = \sum_{q=0}^{\infty} a_{4q+r} \phi_{4q+r}(t)$$

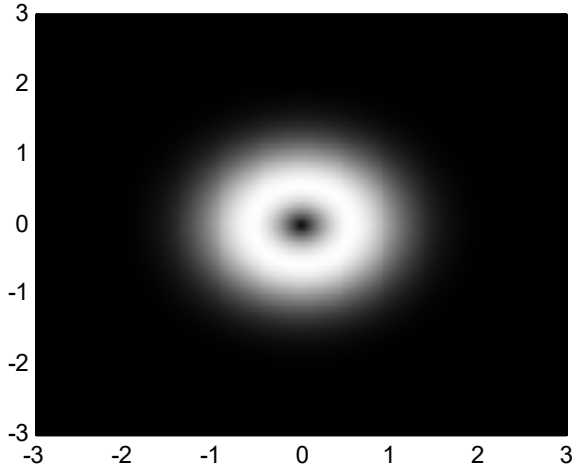
where $r = 0$, 1, 2, or 3,
$a_{4q+r}$ are some constants

$$\int_{-\infty}^{\infty} k(t) e^{-j2\pi f t} dt = (-j)^r k(f)$$

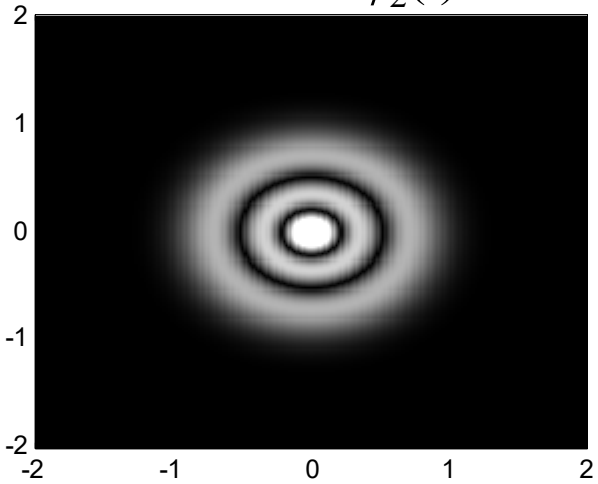
WDF for $\phi_1(t)$



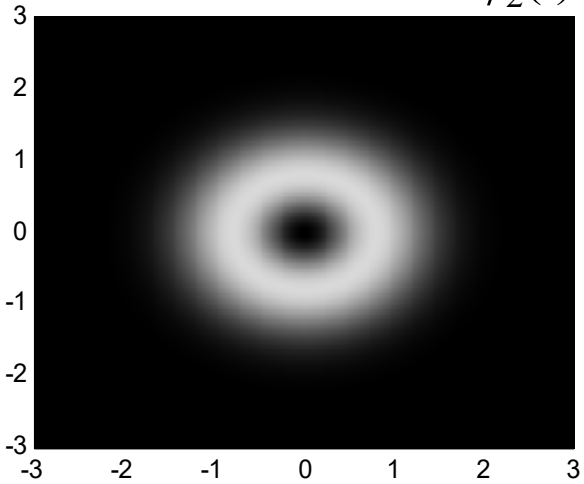

Gabor transform for $\phi_1(t)$




WDF for $\phi_2(t)$




Gabor transform for $\phi_2(t)$



Problem: How to rotate the time-frequency distribution by the angle other than $\pi/2$, $\pi$, and $3\pi/2$?

# 8-3 Rotation: Fractional Fourier Transforms (FRFTs)

$$X_\phi(u) = \sqrt{1 - j\cot\phi}\; e^{j\pi\cot\phi\cdot u^2} \int_{-\infty}^{\infty} e^{-j2\pi\cdot\csc\phi\cdot u\,t}\, e^{j\pi\cdot\cot\phi\cdot t^2}\, x(t)\,dt \;, \qquad \phi = 0.5a\pi$$

When $\phi = 0.5\pi$, the FRFT becomes the FT.

Additivity property:

If we denote the FRFT as $O_F^\phi$ (i.e., $X_\phi(u) = O_F^\phi[x(t)]$)

then $O_F^\sigma\left\{O_F^\phi[x(t)]\right\} = O_F^{\phi+\sigma}[x(t)]$

Physical meaning: Performing the FT *a* times.

Another definition $$X_{\phi}(u)=\sqrt{\frac{1-j\cot\phi}{2\pi}}\,e^{j\frac{\cot\phi}{2}\cdot u^2}\int_{-\infty}^{\infty}e^{-j\csc\phi\cdot u\,t}\,e^{j\frac{\cot\phi}{2}\cdot t^2}\,x(t)dt$$


[Ref] H. M. Ozaktas, Z. Zalevsky, and M. A. Kutay, *The Fractional Fourier Transform with Applications in Optics and Signal Processing*, New York, John Wiley & Sons, 2000.

[Ref] N. Wiener, "Hermitian polynomials and Fourier analysis," *Journal of Mathematics Physics MIT*, vol. 18, pp. 70-73, 1929.

[Ref] V. Namias, "The fractional order Fourier transform and its application to quantum mechanics," *J. Inst. Maths. Applics.*, vol. 25, pp. 241-265, 1980.

[Ref] L. B. Almeida, "The fractional Fourier transform and time-frequency representations," *IEEE Trans. Signal Processing*, vol. 42, no. 11, pp. 3084-3091, Nov. 1994.

[Ref] S. C. Pei and J. J. Ding, "Closed form discrete fractional and affine Fourier transforms," *IEEE Trans. Signal Processing*, vol. 48, no. 5, pp. 1338-1353, May 2000.

$$FT\left[x(t)\right] = X(f)$$

$$FT\left\{FT\left[x(t)\right]\right\} = x(-t)$$

$$FT\left(FT\left\{FT\left[x(t)\right]\right\}\right) = X(-f) = IFT\left[f(t)\right]$$

$$FT\left[FT\left(FT\left\{FT\left[x(t)\right]\right\}\right)\right] = x(t)$$

What happen if we do the FT non-integer times?

**Physical Meaning:**

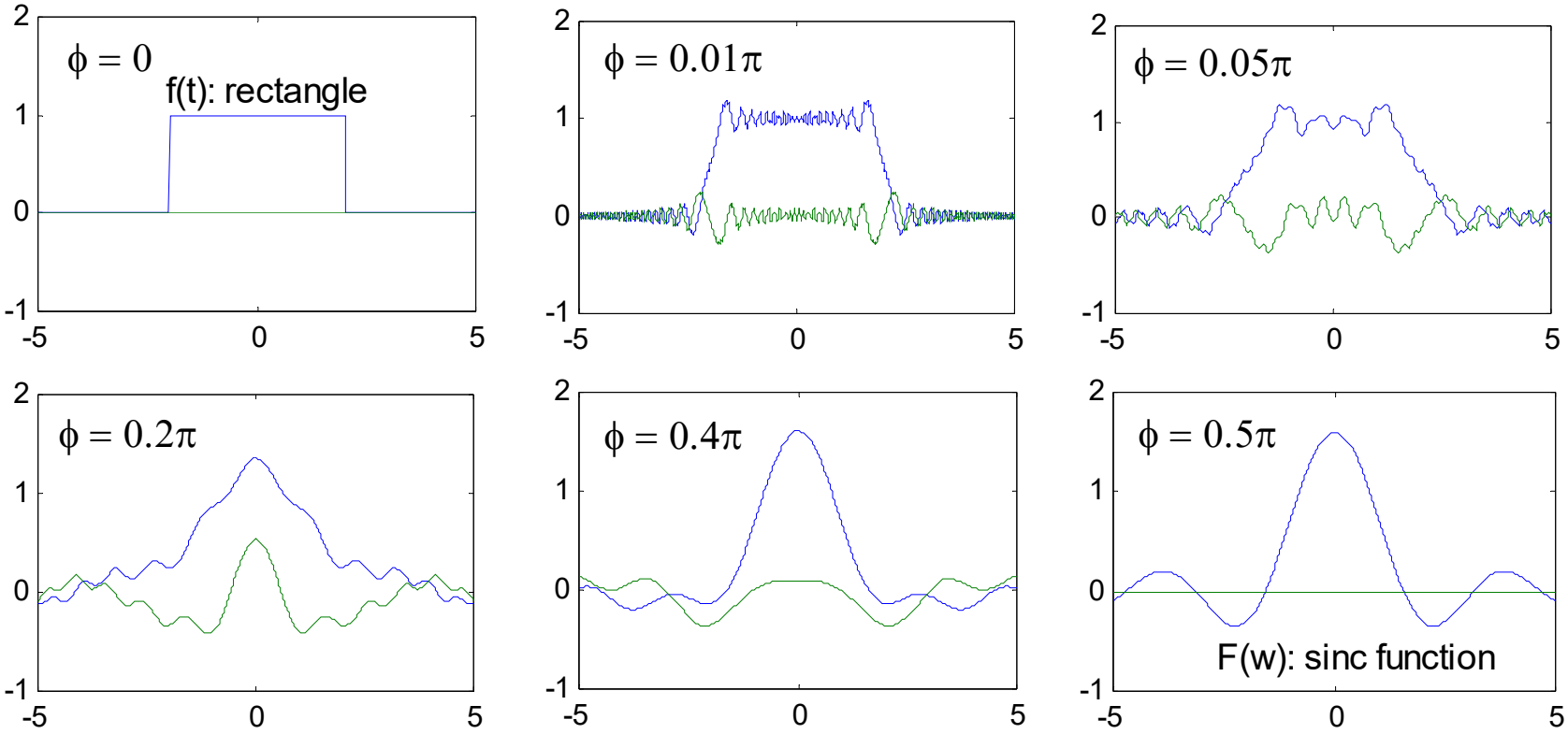
Fourier Transform: time domain $\rightarrow$ frequency domain

Fractional Fourier transform: time domain $\rightarrow$ fractional domain

Fractional domain: the domain between time and frequency

(partially like time and partially like frequency)

Experiment:




blue line: real part

green line: imaginary part

[Ref] L. B. Almeida, "The fractional Fourier transform and time-frequency representations," *IEEE Trans. Signal Processing*, vol. 42, no. 11, pp. 3084-3091, Nov. 1994.

| Time domain | Frequency domain | fractional domain |
|---|---|---|
| Modulation | Shifting | Modulation + Shifting |
| Shifting | Modulation | Modulation + Shifting |
| Differentiation | $\times j2\pi f$ | Differentiation and $\times j2\pi f$ |
| $\times -j2\pi f$ | Differentiation | Differentiation and $\times -j2\pi f$ |

$$x(t-t_0) \xrightarrow{FT} \exp(-j2\pi f t_0) X(f)$$

$$x(t-t_0) \xrightarrow{\text{fractional FT}} \exp(j\varphi - j2\pi u t_0 \sin\phi) X(u - t_0\cos\phi)$$

$$\varphi = \pi t_0^2 \sin\phi\cos\phi$$

$$\frac{d}{dt}x(t) \xrightarrow{FT} j2\pi f\, X(f)$$

$$\frac{d}{dt}x(t) \xrightarrow{\text{fractional FT}} j2\pi u\, X(u)\sin\phi + \frac{d}{du}X(u)\cos\phi$$

**[Theorem]** The fractional Fourier transform (FRFT) with angle $\phi$ is equivalent to the **clockwise rotation operation** with angle $\phi$ for the Wigner distribution function (or for the Gabor transform)

**FRFT with parameter $\phi$ =** **with angle $\phi$**

For the WDF

If $W_x(t, f)$ is the WDF of $x(t)$, and $W_{X\phi}(u, v)$ is the WDF of $X_\phi(u)$, ($X_\phi(u)$ is the FRFT of $x(t)$), then

$$W_{X_\phi}(u,v) = W_x(u\cos\phi - v\sin\phi, u\sin\phi + v\cos\phi)$$

For the Gabor transform (with standard definition)

If $G_x(t, f)$ is the Gabor transform of $x(t)$,

and $G_{X\phi}(u, v)$ is the Gabor transform of $X_\phi(u)$, then

$$G_{X_\phi}(u,v) = e^{j[-2\pi uv\sin^2\phi+\pi(u^2-v^2)\sin(2\phi)/2]} G_x(u\cos\phi - v\sin\phi, u\sin\phi + v\cos\phi)$$

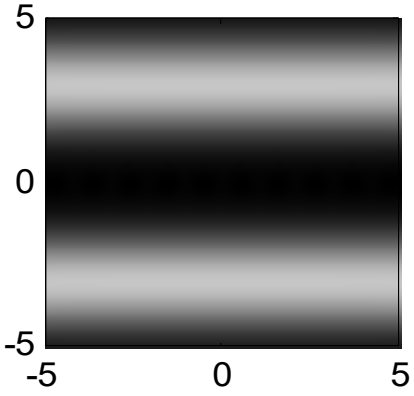
$$\left|G_{X_\phi}(u,v)\right| = \left|G_x(u\cos\phi - v\sin\phi, u\sin\phi + v\cos\phi)\right|$$

For the Gabor transform (with another definition on page 246)

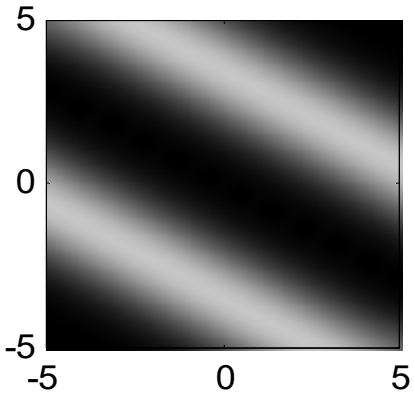
$$G_{X_\phi}(u,v) = G_x(u\cos\phi - v\sin\phi, u\sin\phi + v\cos\phi)$$

The Cohen's class distribution and the Gabor-Wigner transform also have the rotation property

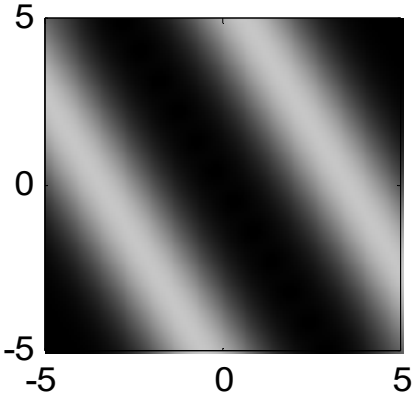
# The Gabor Transform for the FRFT of a cosine function



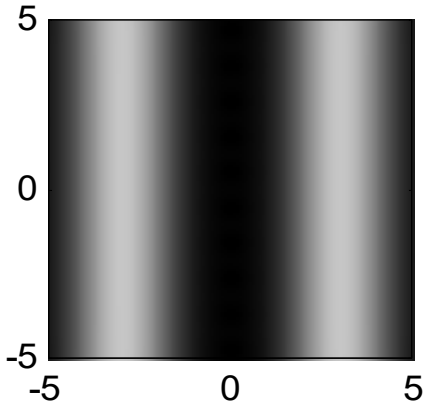

(a) $\phi = 0$



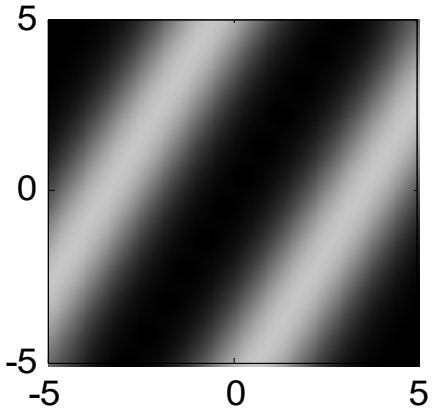

(b) $\phi = \pi/6$



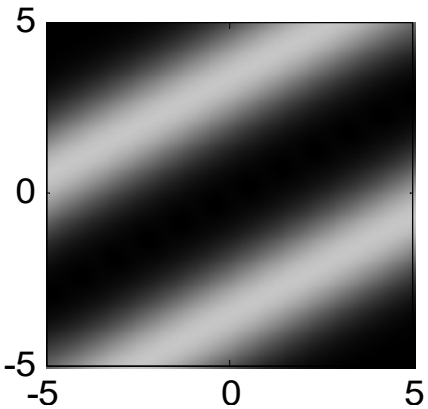

(c) $\phi = 2\pi/6$




(d) $\phi = 3\pi/6$

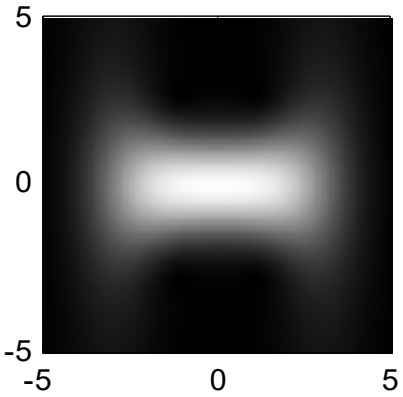



(e) $\phi = 4\pi/6$

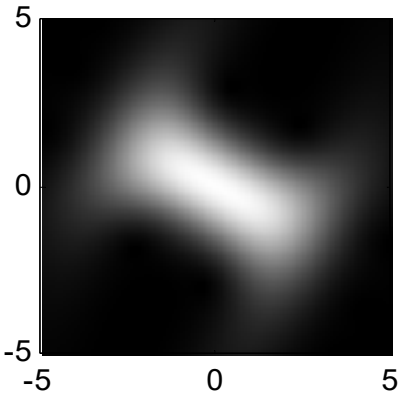



(f) $\phi = 5\pi/6$

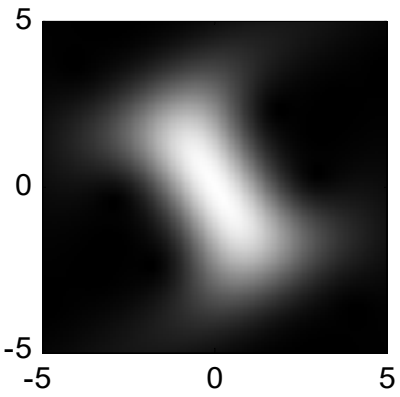
The Gabor Transform for the FRFT of a rectangular function.



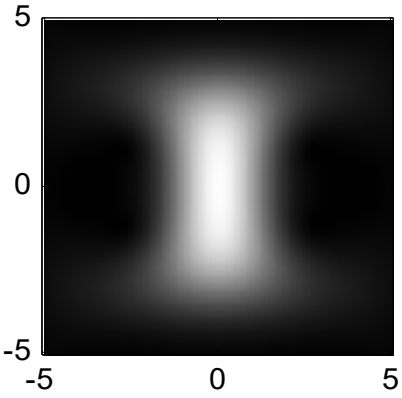

(a) $\phi = 0$



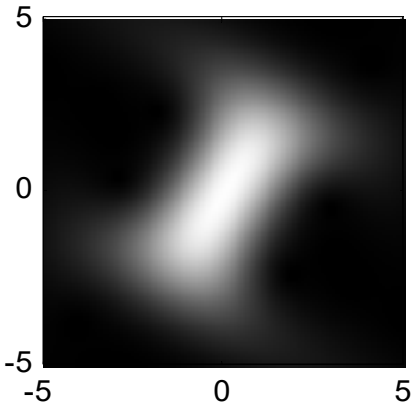

(b) $\phi = \pi/6$



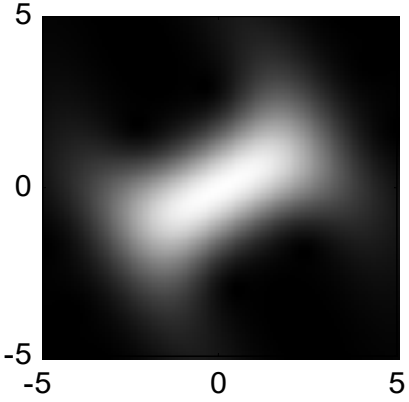

(c) $\phi = 2\pi/6$




(d) $\phi = 3\pi/6$




(e) $\phi = 4\pi/6$




(f) $\phi = 5\pi/6$

# 8-4 Twisting: Linear Canonical Transform (LCT)

$$X_{(a,b,c,d)}(u) = \sqrt{\frac{1}{jb}}\, e^{j\pi \frac{d}{b}u^2} \int_{-\infty}^{\infty} e^{-j2\pi \frac{1}{b}ut}\, e^{j\pi \frac{a}{b}t^2}\, x(t)dt \qquad \text{when } b \neq 0$$

$$X_{(a,0,c,d)}(u) = \sqrt{d}\cdot e^{j\pi cd\,u^2}\, x(d\,u) \qquad \text{when } b = 0$$

$ad - bc = 1$ should be satisfied

Four parameters $a$, $b$, $c$, $d$

## Additivity property of the WDF

If we denote the LCT by $O_F^{(a,b,c,d)}$, i.e., $X_{(a,b,c,d)}(u) = O_F^{(a,b,c,d)}[x(t)]$

then $$O_F^{(a_2,b_2,c_2,d_2)}\left\{O_F^{(a_1,b_1,c_1,d_1)}[x(t)]\right\} = O_F^{(a_3,b_3,c_3,d_3)}[x(t)]$$

where $$\begin{bmatrix} a_3 & b_3 \\ c_3 & d_3 \end{bmatrix} = \begin{bmatrix} a_2 & b_2 \\ c_2 & d_2 \end{bmatrix}\begin{bmatrix} a_1 & b_1 \\ c_1 & d_1 \end{bmatrix}$$

[Ref] K. B. Wolf, "*Integral Transforms in Science and Engineering*," Ch. 9: Canonical transforms, New York, Plenum Press, 1979.

If $W_{X_{(a,b,c,d)}}(u,v)$ is the WDF of $X_{(a,b,c,d)}(u)$, where $X_{(a,b,c,d)}(u)$ is the LCT of $x(t)$, then

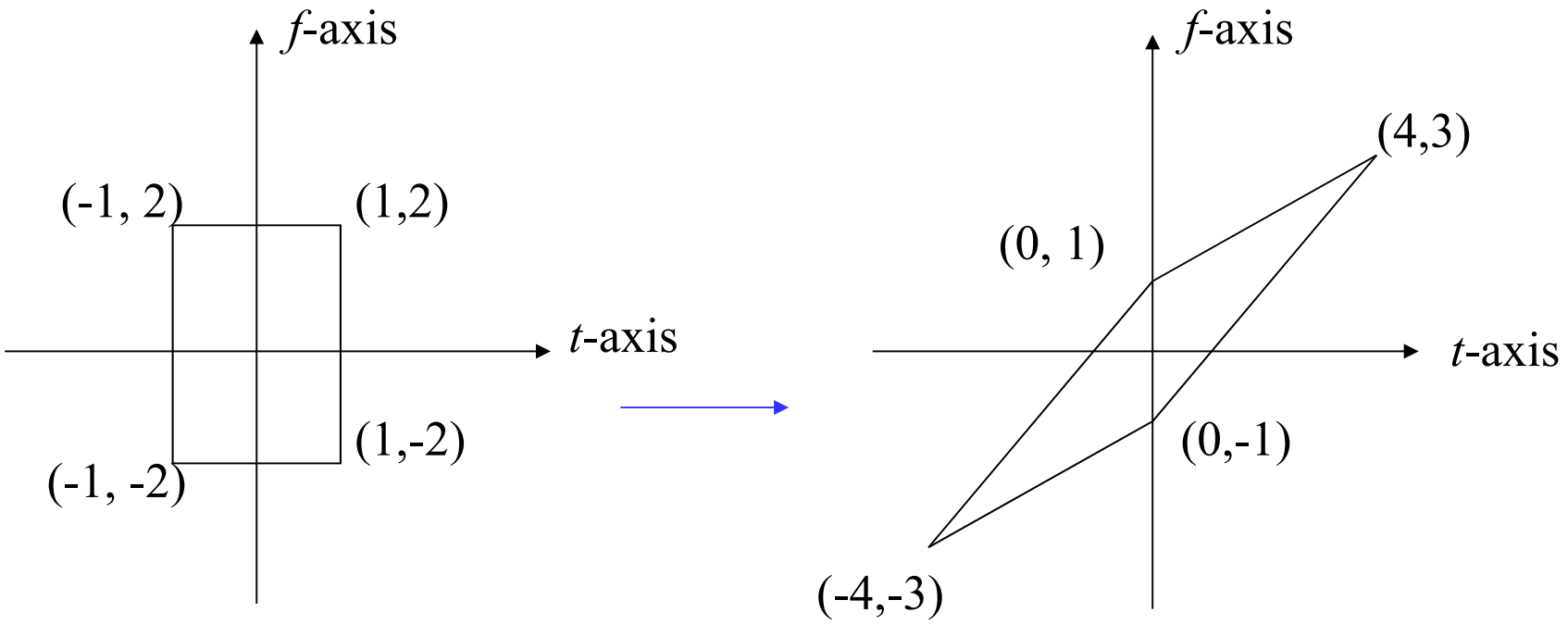
$$W_{X_{(a,b,c,d)}}(u,v) = W_x(du - bv, -cu + av)$$

$$W_{X_{(a,b,c,d)}}(au + bv, cu + dv) = W_x(u,v)$$

LCT == twisting operation for the WDF

The Cohen's class distribution also has the twisting operation.

我們可以自由的用 LCT 將一個中心在 (0, 0) 的平行四邊形的區域，扭曲成另外一個面積一樣且中心也在 (0, 0) 的平行四邊形區域。



$$X_{(a,b,c,d)}(u) = \sqrt{\frac{1}{jb}}\, e^{j\pi\frac{d}{b}u^2} \int_{-\infty}^{\infty} e^{-j2\pi\frac{1}{b}ut}\, e^{j\pi\frac{a}{b}t^2}\, x(t)dt \quad \text{when } b \neq 0$$

$$X_{(a,0,c,d)}(u) = \sqrt{d}\cdot e^{j\pi cdu^2}\, x(du) \quad \text{when } b = 0$$

$ad - bc = 1$ should be satisfied

linear canonical transform

- fractional Fourier transform $\begin{bmatrix} a & b \\ c & d \end{bmatrix} = \begin{bmatrix} \cos\phi & \sin\phi \\ -\sin\phi & \cos\phi \end{bmatrix}$
  - $\phi = \pi/2$ Fourier transform
  - $\phi = 0$ identity operation
  - $\phi = -\pi/2$ inverse Fourier transform
- Fresnel transform (convolution with a chirp) $\begin{bmatrix} a & b \\ c & d \end{bmatrix} = \begin{bmatrix} 1 & \lambda z \\ 0 & 1 \end{bmatrix}$
- chirp multiplication $\begin{bmatrix} a & b \\ c & d \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ \tau & 1 \end{bmatrix}$

$$X_{(a,0,c,d)}(u) = e^{j\pi\tau u^2}\, x(u)$$

- scaling $\begin{bmatrix} a & b \\ c & d \end{bmatrix} = \begin{bmatrix} 1/\sigma & 0 \\ 0 & \sigma \end{bmatrix}$

# 附錄十一　Linear Canonical Transform 和光學系統的關係

(1) Fresnel Transform (電磁波在空氣中的傳播)

$$U_o(x,y) = -\frac{i}{\lambda}\frac{e^{ikz}}{z}\int_{-\infty}^{\infty}\int_{-\infty}^{\infty} e^{j\frac{k}{2z}\left[(x-x_i)^2+(y-y_i)^2\right]} U_i(x_i,y_i)\,dx_i dy_i$$

$k = 2\pi/\lambda$: wave number　　$\lambda$: wavelength　　$z$: distance of propagation

$$U_o(x,y) = e^{ikz}\sqrt{\frac{1}{j\lambda z}}\int_{-\infty}^{\infty} e^{j\frac{k}{2z}(y-y_i)^2}\sqrt{\frac{1}{j\lambda z}}\int_{-\infty}^{\infty} e^{j\frac{k}{2z}(x-x_i)^2} U_i(x_i,y_i)\,dx_i dy_i$$

(2 個 1-D 的 LCT)

$$e^{j\frac{k}{2z}x^2} * U_i(x,y_i) = e^{j\frac{\pi}{\lambda z}x^2} * U_i(x,y_i)$$

Fresnel transform 相當於 LCT $\begin{bmatrix} a & b \\ c & d \end{bmatrix} = \begin{bmatrix} 1 & \lambda z \\ 0 & 1 \end{bmatrix}$

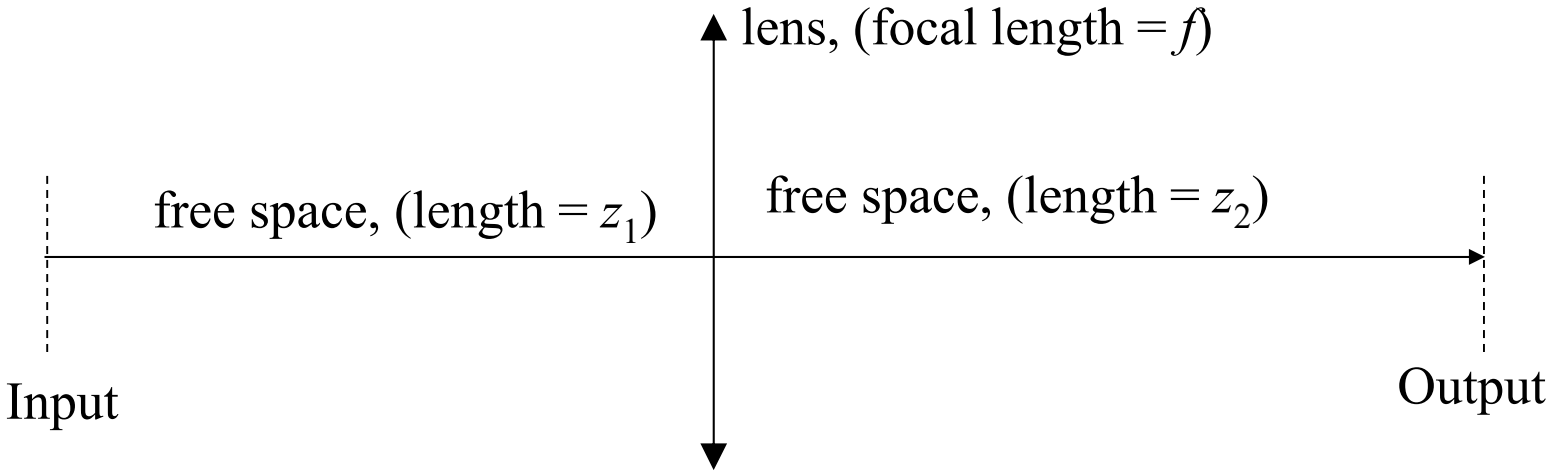
(2) Spherical lens, refractive index = $n$

$$U_o(x,y) = e^{ikn\Delta} e^{-j\frac{k}{2f}\left[x^2+y^2\right]} U_i(x,y)$$

$f$ : focal length  $\Delta$: thickness of lens

經過 lens 相當於 LCT $\begin{bmatrix} a & b \\ c & d \end{bmatrix} = \begin{bmatrix} 1 & 0 \\ -1/\lambda f & 1 \end{bmatrix}$ 的情形

(3) Free spaces + Spherical lens




Input 和 output 之間的關係，可以用 LCT 表示

$$\begin{bmatrix} a & b \\ c & d \end{bmatrix} = \begin{bmatrix} 1 & \lambda z_2 \\ 0 & 1 \end{bmatrix} \begin{bmatrix} 1 & 0 \\ -1/\lambda f & 1 \end{bmatrix} \begin{bmatrix} 1 & \lambda z_1 \\ 0 & 1 \end{bmatrix} = \begin{bmatrix} 1-\frac{z_2}{f} & \lambda(z_1+z_2)-\frac{\lambda z_1 z_2}{f} \\ -\frac{1}{\lambda f} & 1-\frac{z_1}{f} \end{bmatrix}$$

LCT: $X_{(a,b,c,d)}(u)=\sqrt{\frac{1}{jb}}\,e^{j\pi\frac{d}{b}u^2}\int_{-\infty}^{\infty}e^{-j2\pi\frac{1}{b}ut}e^{j\pi\frac{a}{b}t^2}\,x(t)dt$

$$\begin{bmatrix}a & b\\ c & d\end{bmatrix}=\begin{bmatrix}1-\frac{z_2}{f} & \lambda(z_1+z_2)-\frac{\lambda z_1 z_2}{f}\\ -\frac{1}{\lambda f} & 1-\frac{z_1}{f}\end{bmatrix}$$

$z_1 = z_2 = 2f \rightarrow$ 即高中物理所學的「倒立成像」

$$\begin{bmatrix}a & b\\ c & d\end{bmatrix}=\begin{bmatrix}-1 & 0\\ -\frac{1}{\lambda f} & -1\end{bmatrix}$$

$z_1 = z_2 = f \rightarrow$ Fourier Transform + Scaling

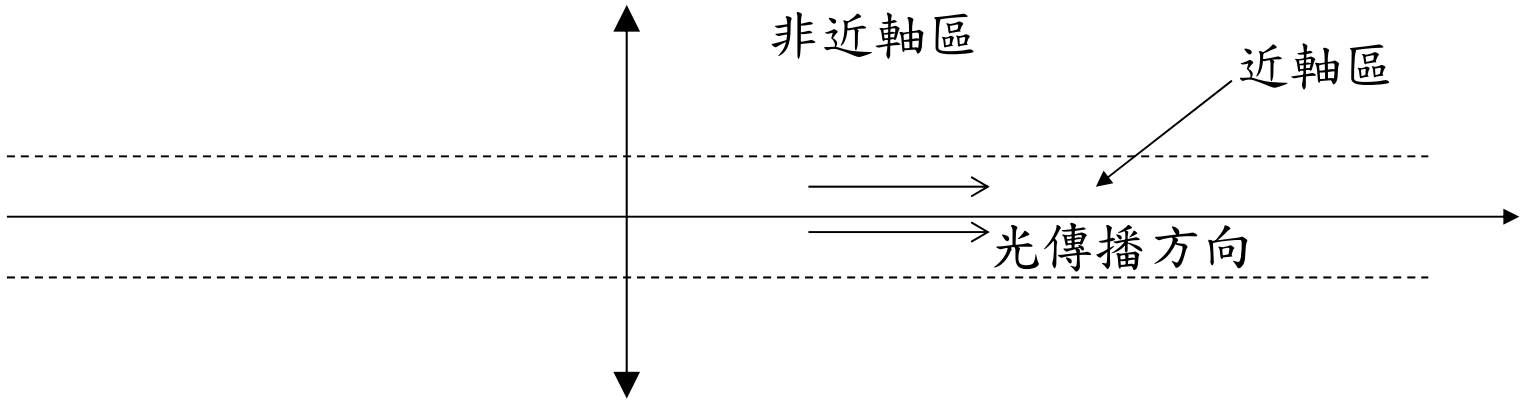
$$\begin{bmatrix}a & b\\ c & d\end{bmatrix}=\begin{bmatrix}0 & \lambda f\\ -\frac{1}{\lambda f} & 0\end{bmatrix}$$

$z_1 = z_2 \rightarrow$ fractional Fourier Transform + Scaling

用 LCT 來分析光學系統的好處：

只需要用到 2×2 的矩陣運算，避免了複雜的物理理論和數學積分

但是 LCT 來分析光學系統的結果，只有在「近軸」的情形下才準確




參考資料：

[1] H. M. Ozaktas and D. Mendlovic, “Fractional Fourier optics,” *J. Opt. Soc. Am. A*, vol.12, 743-751, 1995.

[2] L. M. Bernardo, “ABCD matrix formalism of fractional Fourier optics,” *Optical Eng.*, vol. 35, no. 3, pp. 732-740, March 1996.